

Photonic phase-gate using Rydberg atoms and Microwaves

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Aim

Design a quasi-deterministic phase gate for photonic qubits using cold Rydberg atoms

EIT + Rydbergs

allow storage of single excitations in an atomic cloud and provide strong nonlinear interactions [1-4,7]

Separate storage and interactions

using microwaves and an auxiliary state to prevent loss.

Nonlocal character

of the blockade might allow circumvention of conventional limitations to local (Kerr-like) nonlinearities [7-9]

dual rail photonic qubits

EIT

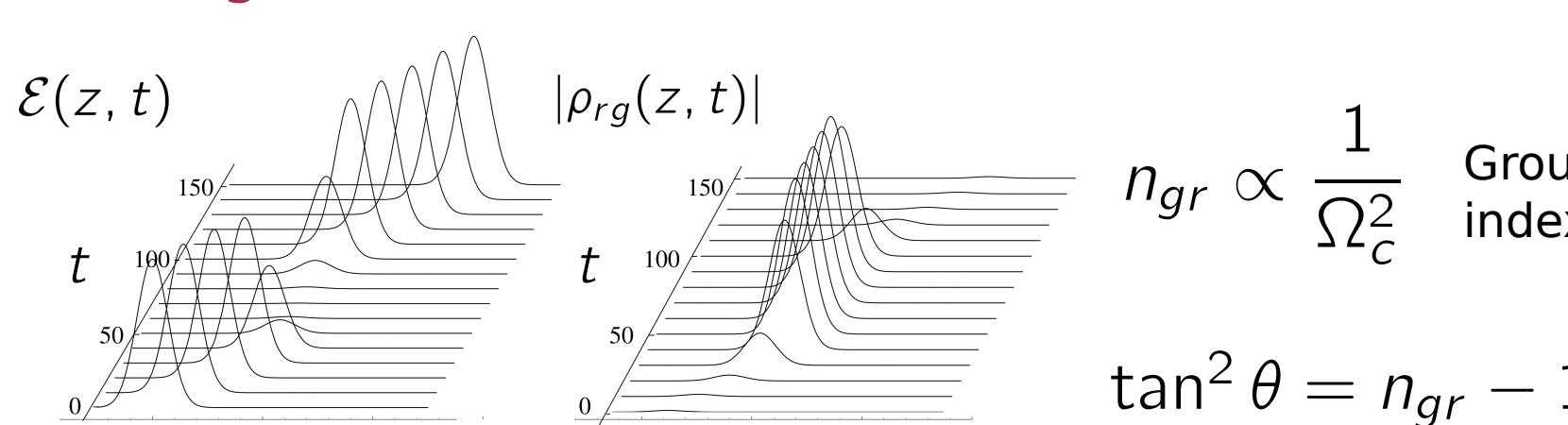
atomic excitations (polaritons)

Using **electromagnetically induced transparency (EIT)**, photonic qubits are coherently mapped into a spin wave (polariton) [14,15] with Rydberg content.

$$\Psi(z, t) = \cos \theta \mathcal{E}(z, t) - \sin \theta \sqrt{N} \rho_{rg}(z, t) e^{i \Delta k z}$$

light

atomic excitation



Drive a 2π rotation in the target qubit through an auxiliary state

$$|p\rangle = \sum_i^N e^{i\phi_i} |p_i\rangle \quad |p_i\rangle = |g\dots p_i\dots g\rangle$$

In the non-interacting case, it adds a global (boring) π phase shift to the wavefunction.

However, microwaves turn on resonant dipole-dipole interactions between stored polaritons [4].

$$R_b^{(6)} = \sqrt[6]{C_3 / \Omega_\mu} \quad \text{Coherence length of vDW interaction}$$

$$R_b^{(3)} = \sqrt[3]{C_3 / \Omega_\mu} \quad \text{Microwave-induced}$$

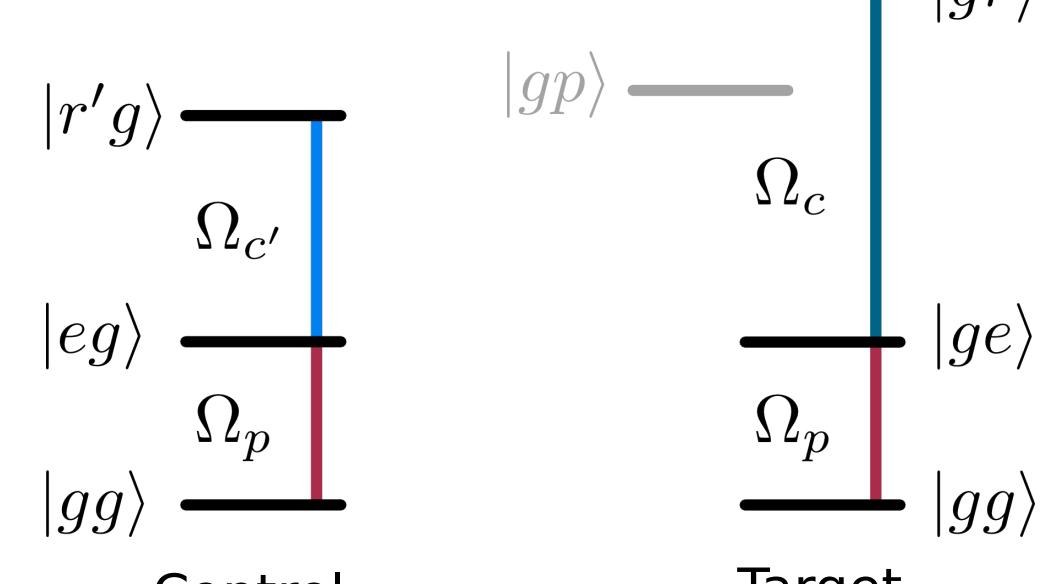
$$n_{gr} \propto \frac{1}{\Omega_c^2} \quad \text{Group index}$$

$$\tan^2 \theta = n_{gr} - 1$$

Polaritons involve **Rydberg levels with strong interactions**, which stores each photon as a **collective excitation**.

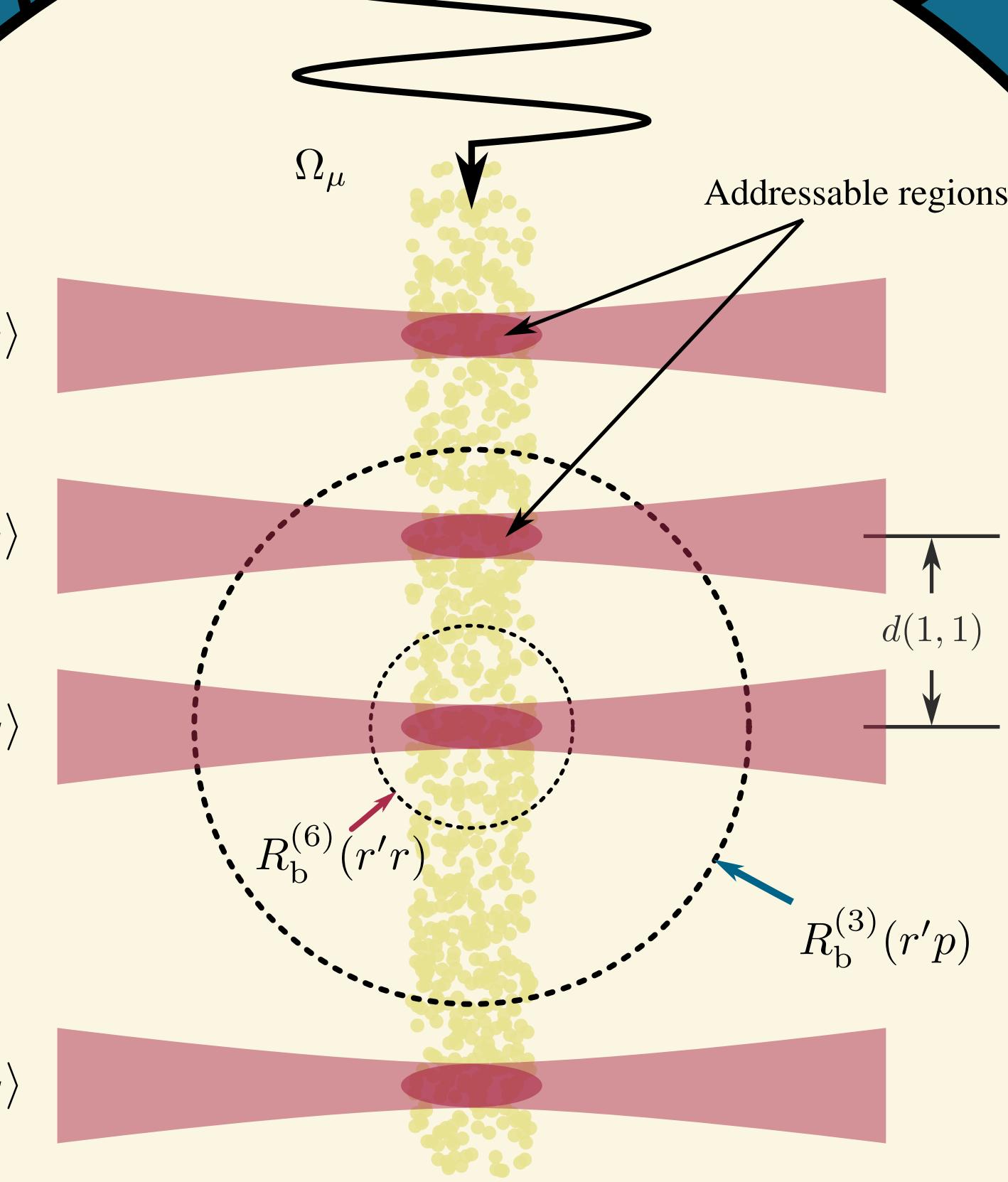
$$|r\rangle = \sum_i^N e^{i\phi_i} |r_i\rangle \quad \text{The summation is carried over states with a single excitation}$$

$|r_i\rangle = |g\dots r_i\dots g\rangle$



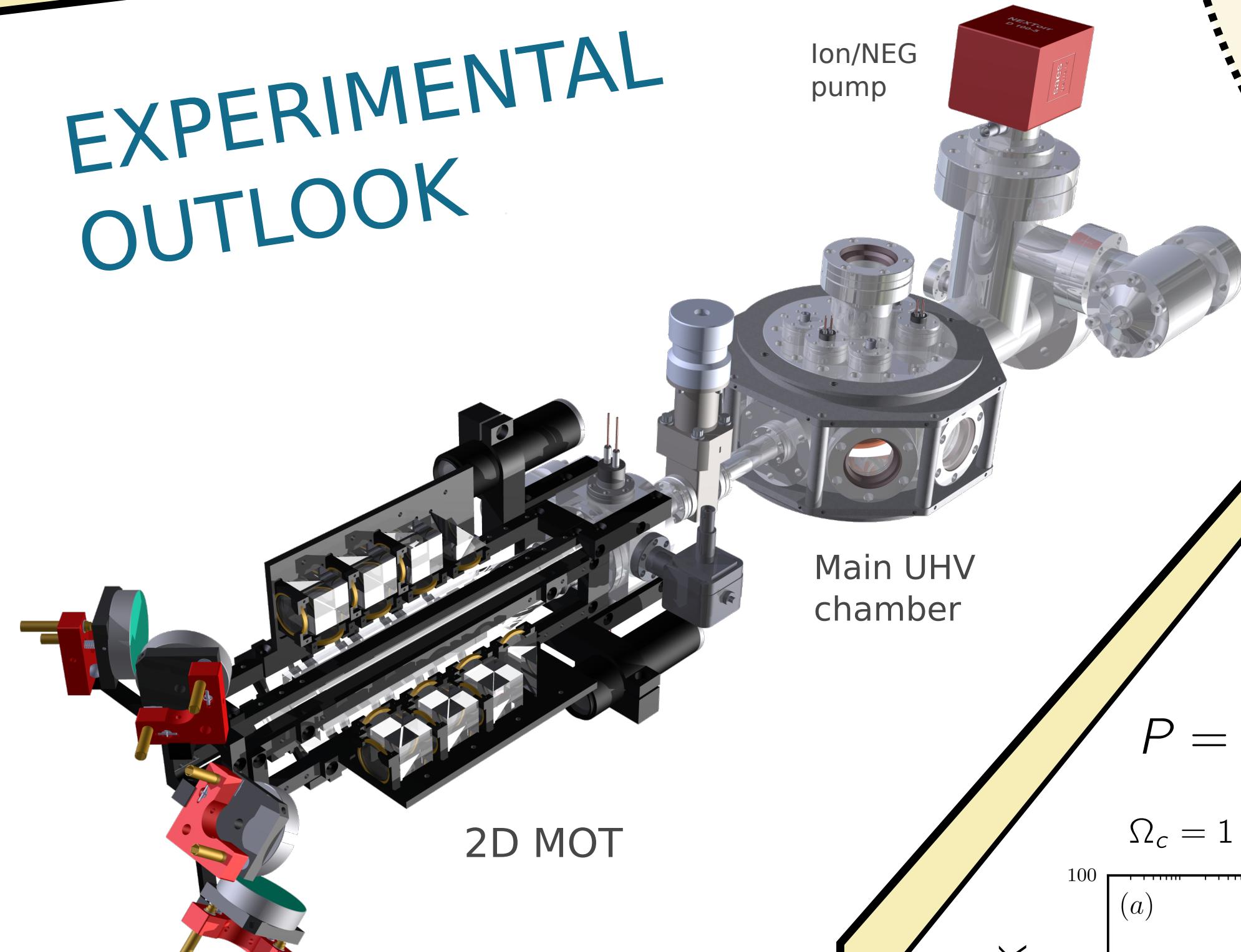
Noninteracting level system for the control (left) and target (right) states in the inner region.

MICROWAVE DRIVING

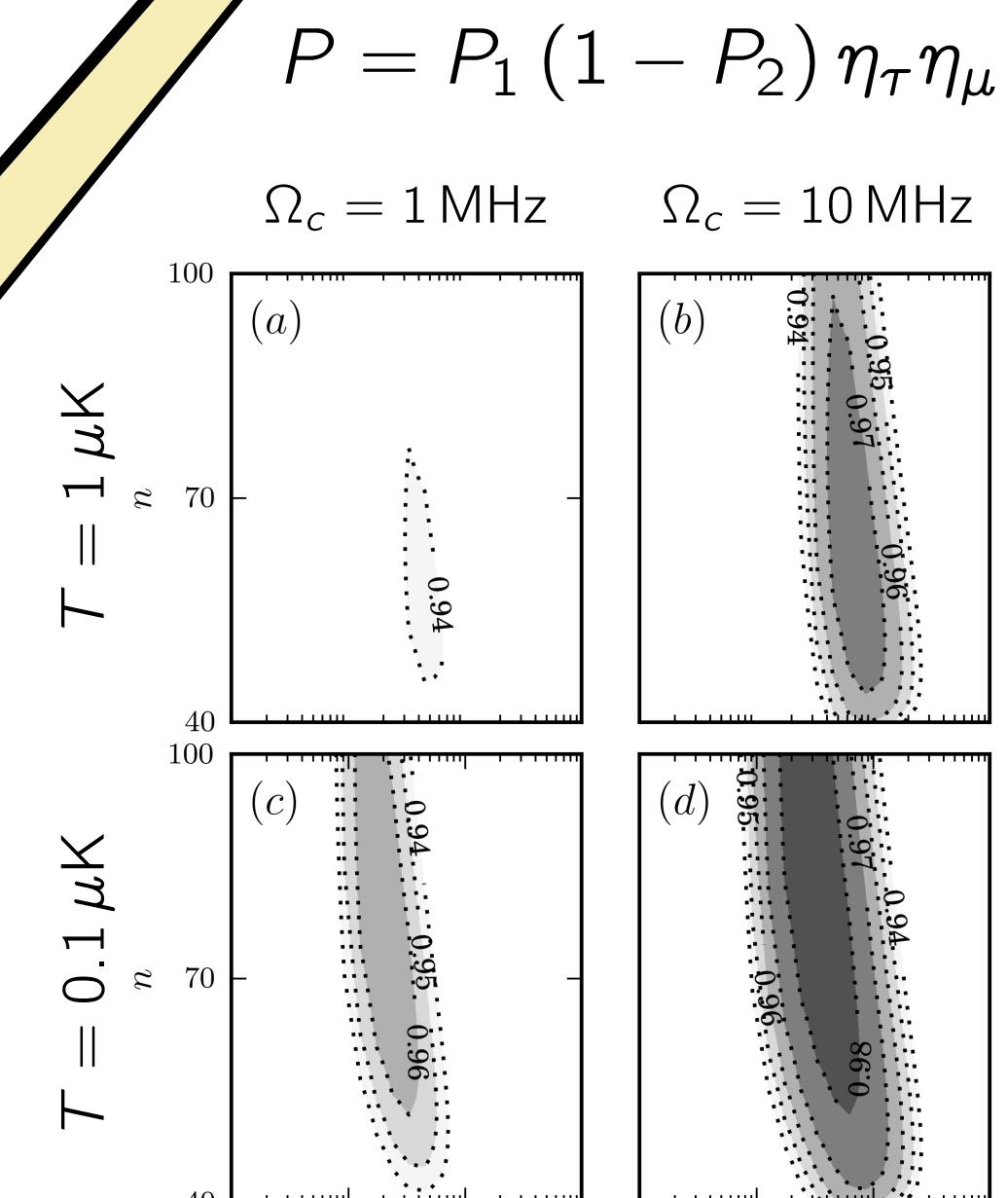


ESTIMATION

EXPERIMENTAL OUTLOOK



- Higher repetition rates thanks to efficient 2D MOT loading
- Higher magnetic field gradients for compressed MOT phase to increase atomic density in dipole trap.
- In-vacuum high NA lenses to address individual polaritons with tightly focused beams
- Built-in microwave control



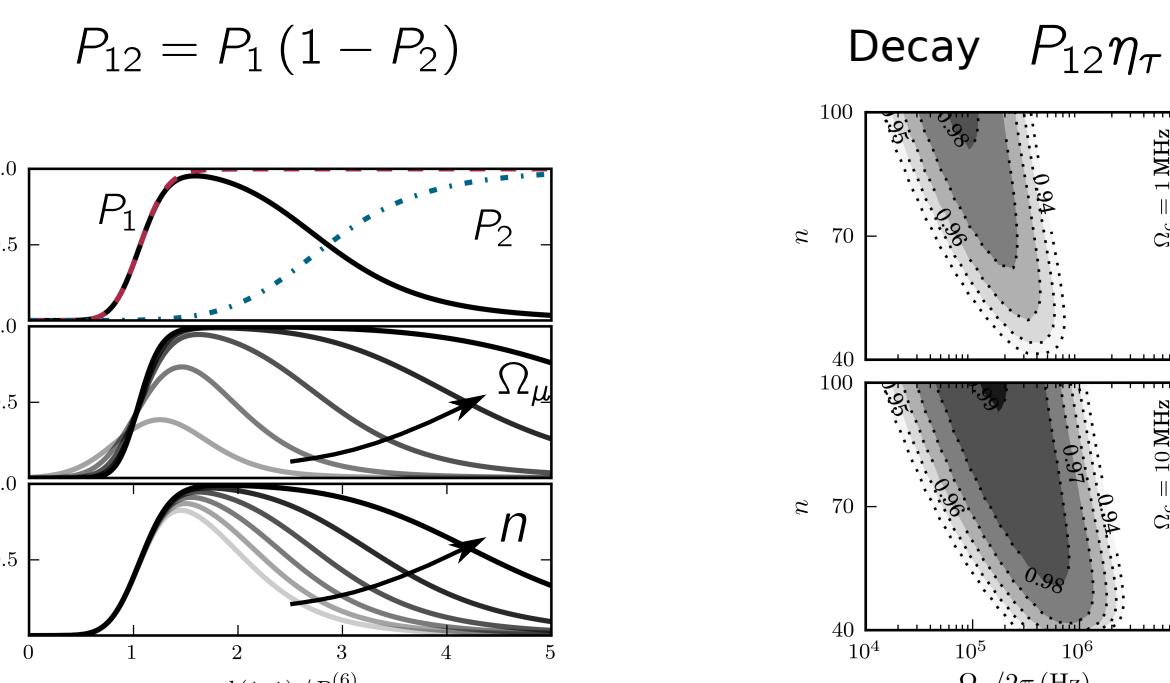
Look at the success probability of two phenomena in the interaction region:

P_1 - population of the state $|1_C, 1_T\rangle = |r'r\rangle$ given a Rydberg excitation in the control site

P_2 - probability to drive a 2π transition through $|1_C, 1_T\rangle = |r'p\rangle$

The probability of success is then given by $P_{12} = P_1 (1 - P_2)$

$$P_{12} = P_1 (1 - P_2)$$



Motional dephasing [16]

$$\eta_m = \frac{e^{-\epsilon t^2/\zeta^2}}{(1+t^2/\zeta^2)^2} e^{-t^2/\zeta^2}$$

ξ Atoms moving in/out of the region

ζ Atoms moving within the spin-wave grating



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